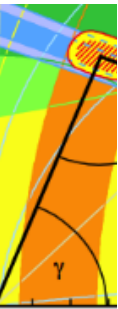


# Messung von $\gamma$ in Tree-Zerfällen bei LHCb

Till Moritz Karbach  
CERN

moritz.karbach @ cern.ch

DPG, Dresden, March 2013

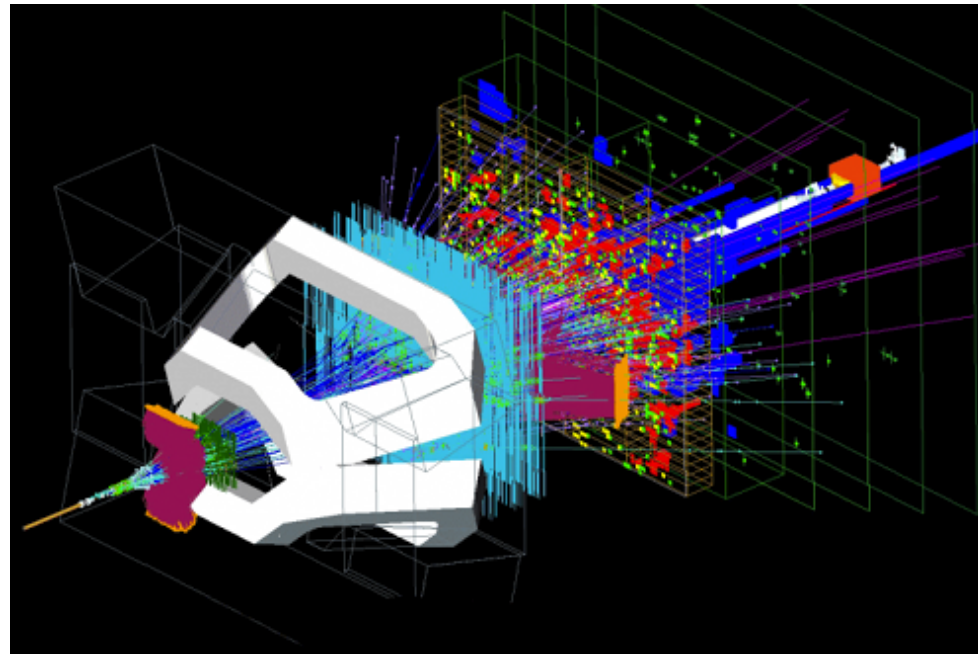
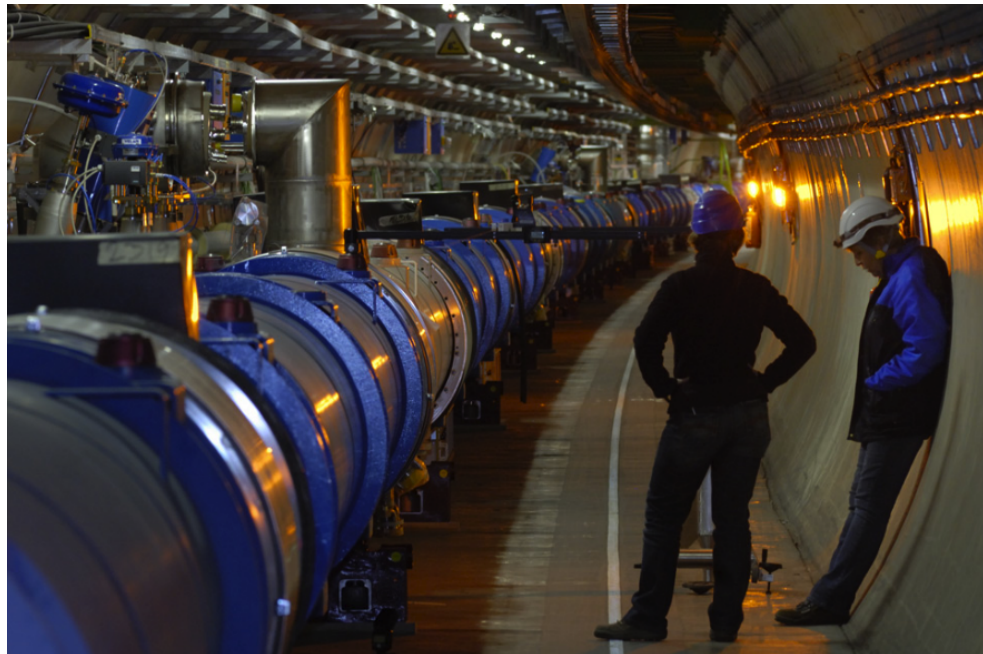


# Plan

- Part I: introduction
- Part II: time *integrated*  $\gamma$  measurements
- Part III: combination of time int. meas.
- Part IV: time *dependent*  $\gamma$  measurement

# LHCb

- LHCb is a forward spectrometer at LHC.
- Operating in collider mode.
- It is dedicated to precision measurements of b and c physics ...  
... a general purpose spectrometer in the forward region!
- CP violation, rare decays



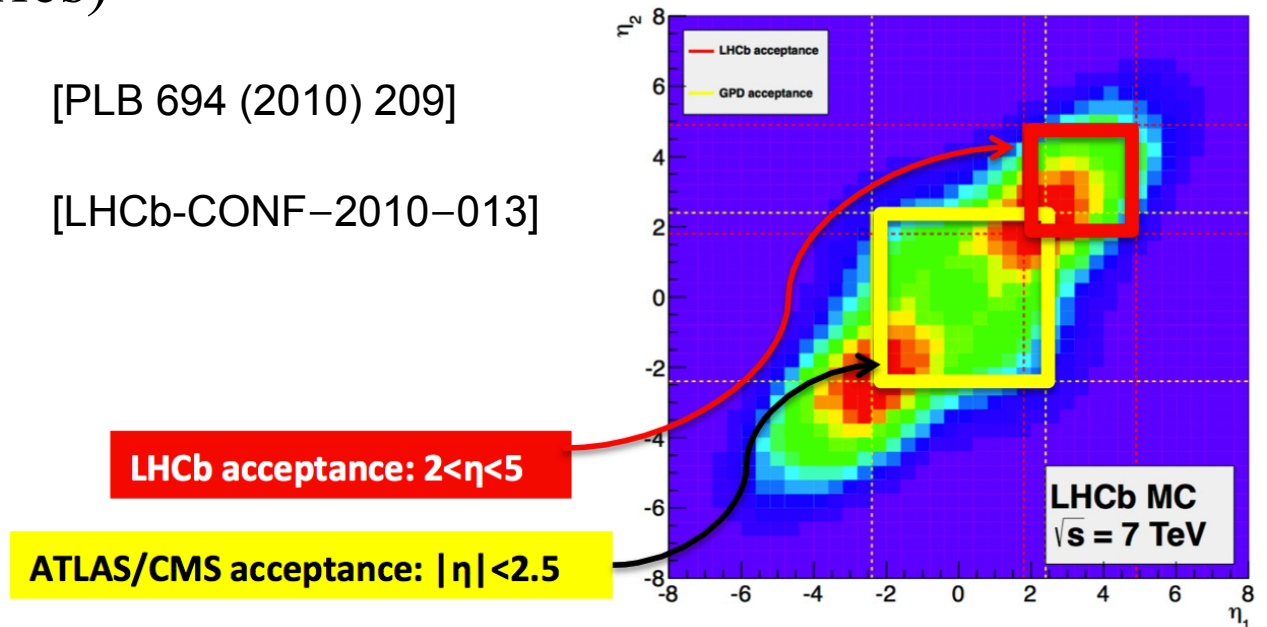
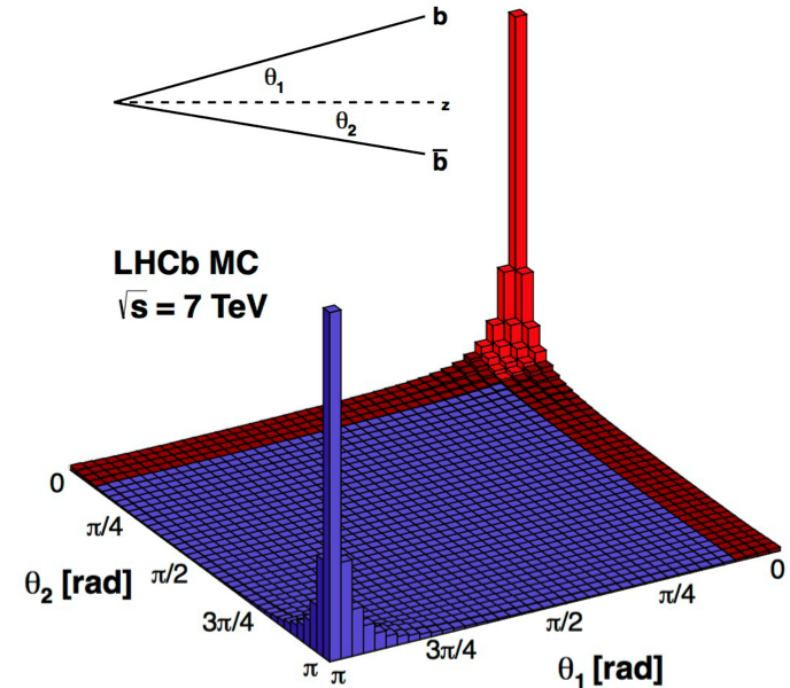
# LHCb

- one arm forward spectrometer
- $b$  pair production angles strongly correlated
- covers  $1.9 < \eta < 4.9$
- 100'000  $b\bar{b}$  pairs produced per second ( $10^4 \times B$  factories)

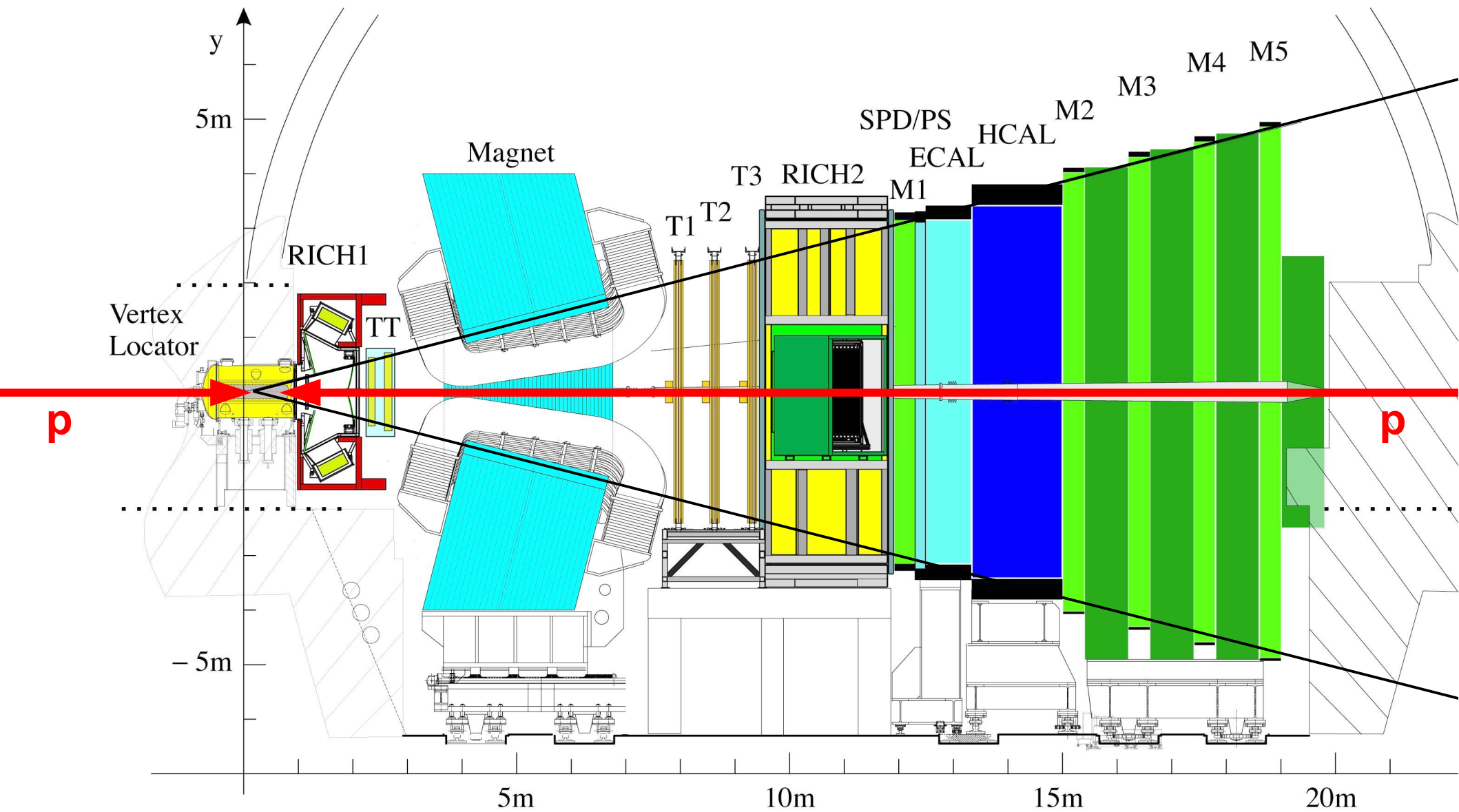
$$\sigma(b\bar{b}) = 284 \pm 53 \mu\text{b} \quad [\text{PLB } 694 (2010) 209]$$

$$\sigma(c\bar{c}) \approx 20 \times \sigma(b\bar{b}) \quad [\text{LHCb-CONF-2010-013}]$$

- particle identification by two RICH detectors



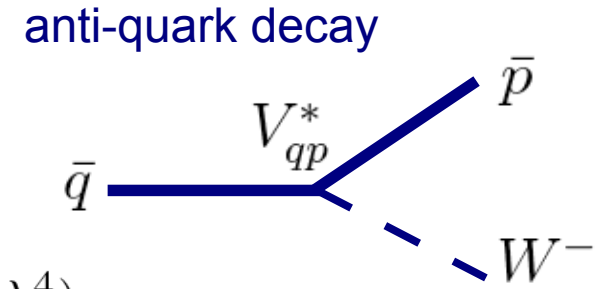
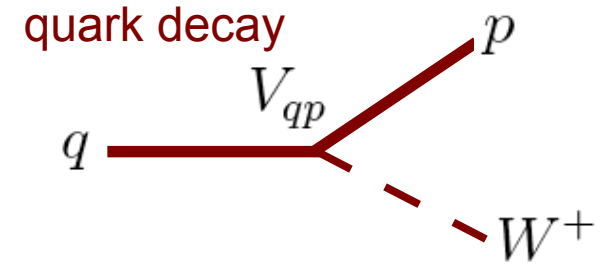
# LHCb



# CKM matrix

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \cdot \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

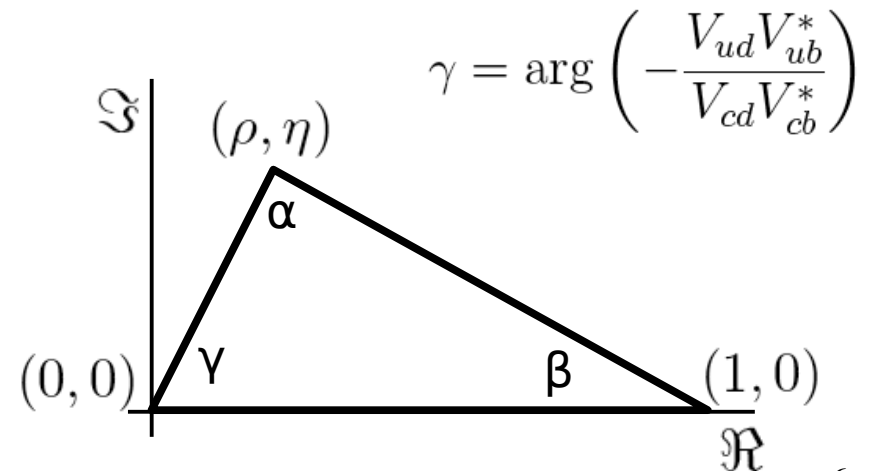
flavor eigenstates mass eigenstates



$$V_{\text{CKM}} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & \lambda^3 A(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & \lambda^2 A \\ \lambda^3 A(1 - \rho - i\eta) & -\lambda^2 A & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

unitarity condition:  $V^\dagger V = 1$

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$



# angle $\gamma$

$\gamma$  is the least well known angle of the unitarity triangle.

“combined  $\gamma$  measurements”

$$\gamma = (66^{+12}_{-12})^\circ$$

$$\gamma = (75.5 \pm 10.5)^\circ$$

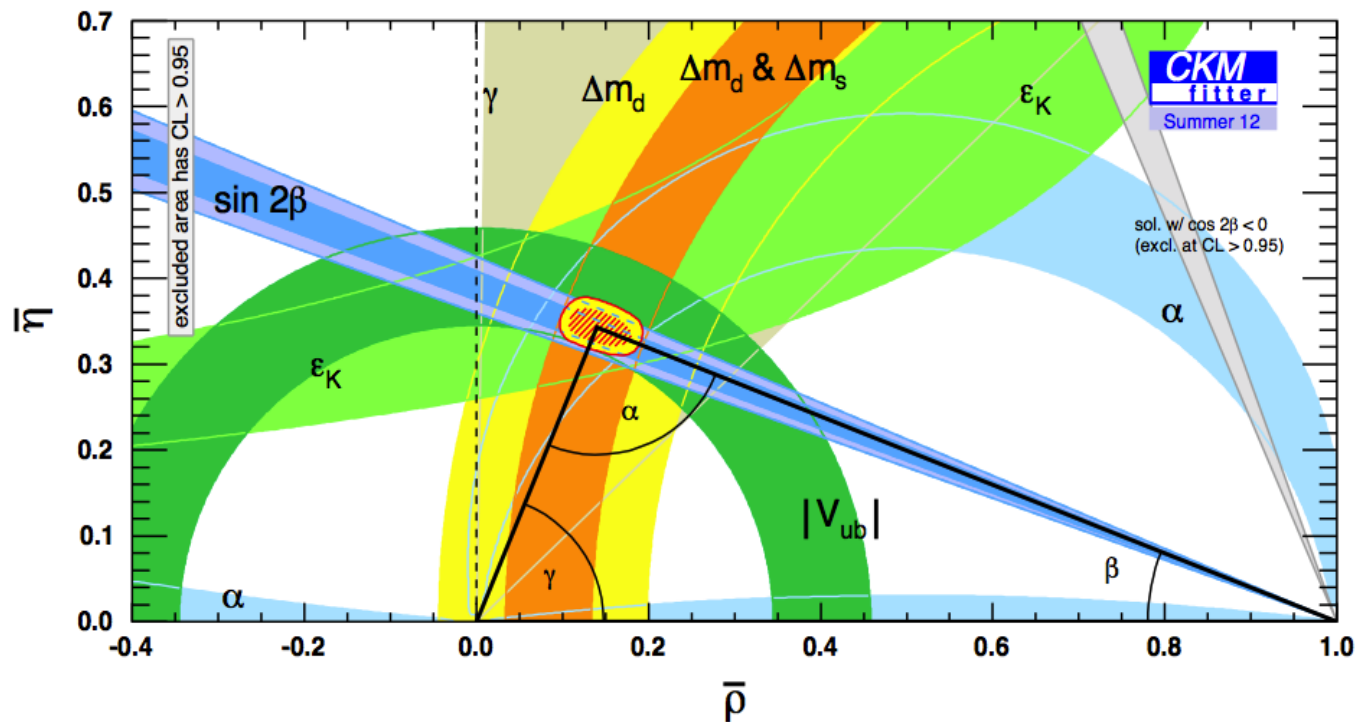
CKMfitter ICHEP 2012

UTfit pre-ICHEP 2012

“ $\gamma$  meas. not in triangle fit”

$$\gamma = (68.0^{+4.1}_{-4.6})^\circ$$

$$\gamma = (67.8 \pm 3.2)^\circ$$



# angle $\gamma$

- Difficult to measure, as the decay rates are small.

$$\text{BR}(B^- \rightarrow DK^-, D \rightarrow \pi K) \approx 2 \times 10^{-7} \quad \text{LHCb first observation with } \sim 100 \text{ events}$$

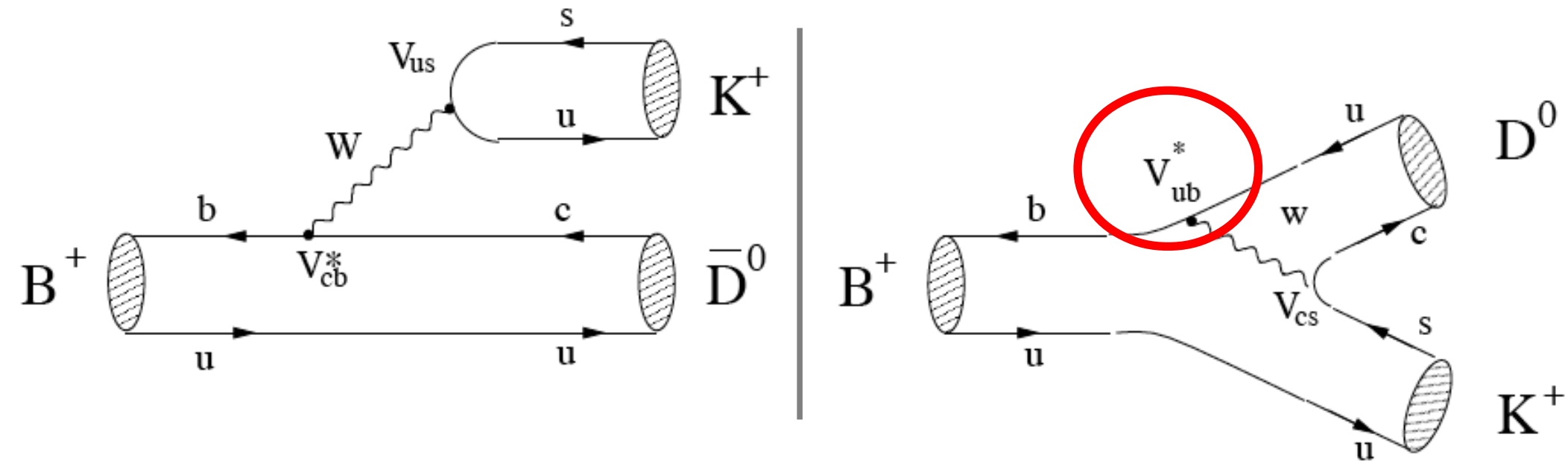
- $\gamma$  can be determined entirely from **tree decays**.
  - this is a **unique** property among all CP violation parameters
  - **negligible** theoretical uncertainty (J. Zupan):  $\delta\gamma/\gamma = \mathcal{O}(10^{-6})$
  - hadronic parameters can all be **determined from the data**
  - provides an important **Standard Model set point** (“standard candle”)
- $\gamma$  can probe for new physics at extremely **high energy scales** (J. Zupan)
  - (N)MFV new physics scenarios:  $\sim \mathcal{O}(10^2 \text{ TeV})$
  - gen. FV new physics scenarios:  $\sim \mathcal{O}(10^3 \text{ TeV})$



Part II:

time integrated measurements

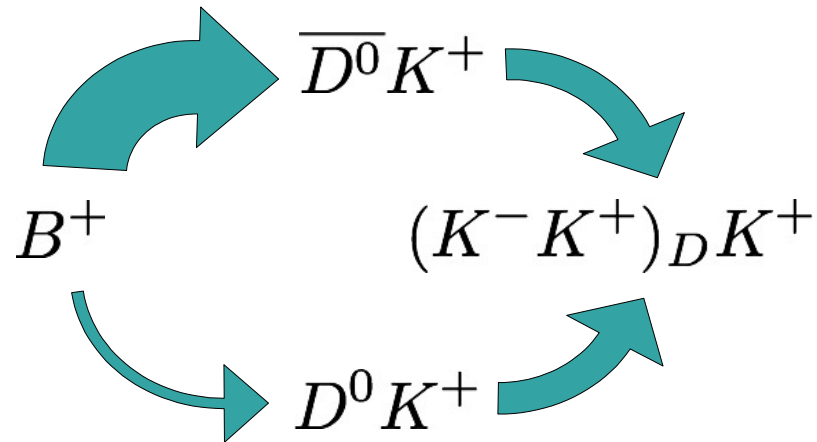
# $B \rightarrow DK$



- This was, and still is, the most important channel to measure  $\gamma$ .
- We need to reconstruct the  $D/\bar{D}$  meson in a final state accessible to both to achieve interference.
- **Also possible:  $B \rightarrow D\pi$ !** But little sensitivity.
- Choice of final state labels the “method”: GLW, ADS, GGSZ

# B $\rightarrow$ DK

“GLW”

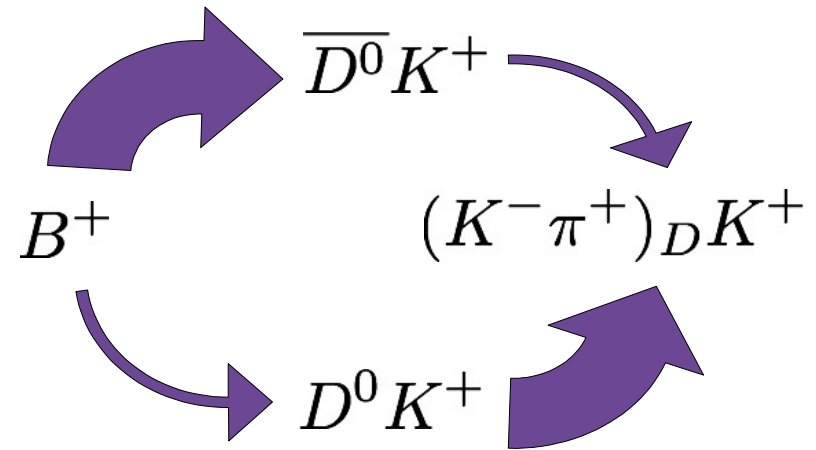


Phys.Lett. B253 (1991) 483

Phys.Lett. B265 (1991) 172

Gronau, London, Wyler

“ADS”, “suppressed”



Phys.Rev.Lett 78 (1997) 3257

Phys.Rev. D63 (2001) 036005

Atwood, Dunietz, Soni

“GGSZ”, “Dalitz”

- Use 3-body self-conjugate modes such as  $D \rightarrow K_S \pi^+ \pi^-$
- hadronic D parameters vary across Dalitz plot

Giri, Grossman, Soffer, Zupan, hep-ph/0303187

# B → Dh: Observables

- Define observables as **yield ratios** (many systematics cancel).
- Charge **asymmetries**:

$$A_h^f = \frac{\Gamma(B^- \rightarrow [f]_D h^-) - \Gamma(B^+ \rightarrow [f]_D h^+)}{\Gamma(B^- \rightarrow [f]_D h^-) + \Gamma(B^+ \rightarrow [f]_D h^+)}$$

- **Kaon/pion** ratio:

$$R_{K/\pi}^f = \frac{\Gamma(B^\pm \rightarrow [f]_D K^\pm)}{\Gamma(B^\pm \rightarrow [f]_D \pi^\pm)}$$

Form a system of equations.  
Need more observables than  
parameters!  
→ many different D decays

- **Suppressed/favored** decay ratio (2-body example):

$$R_h^\pm = \frac{\Gamma(B^\pm \rightarrow [\pi^\pm K^\mp]_D h^\pm)}{\Gamma(B^\pm \rightarrow [K^\pm \pi^\mp]_D h^\pm)}$$

$$= r_B^2 + r_D^2 + 2r_B r_D \cos(\underbrace{\pm\gamma + \delta_B + \delta_D}_{\text{strong phase diff.}})$$

# B → Dh: Observables

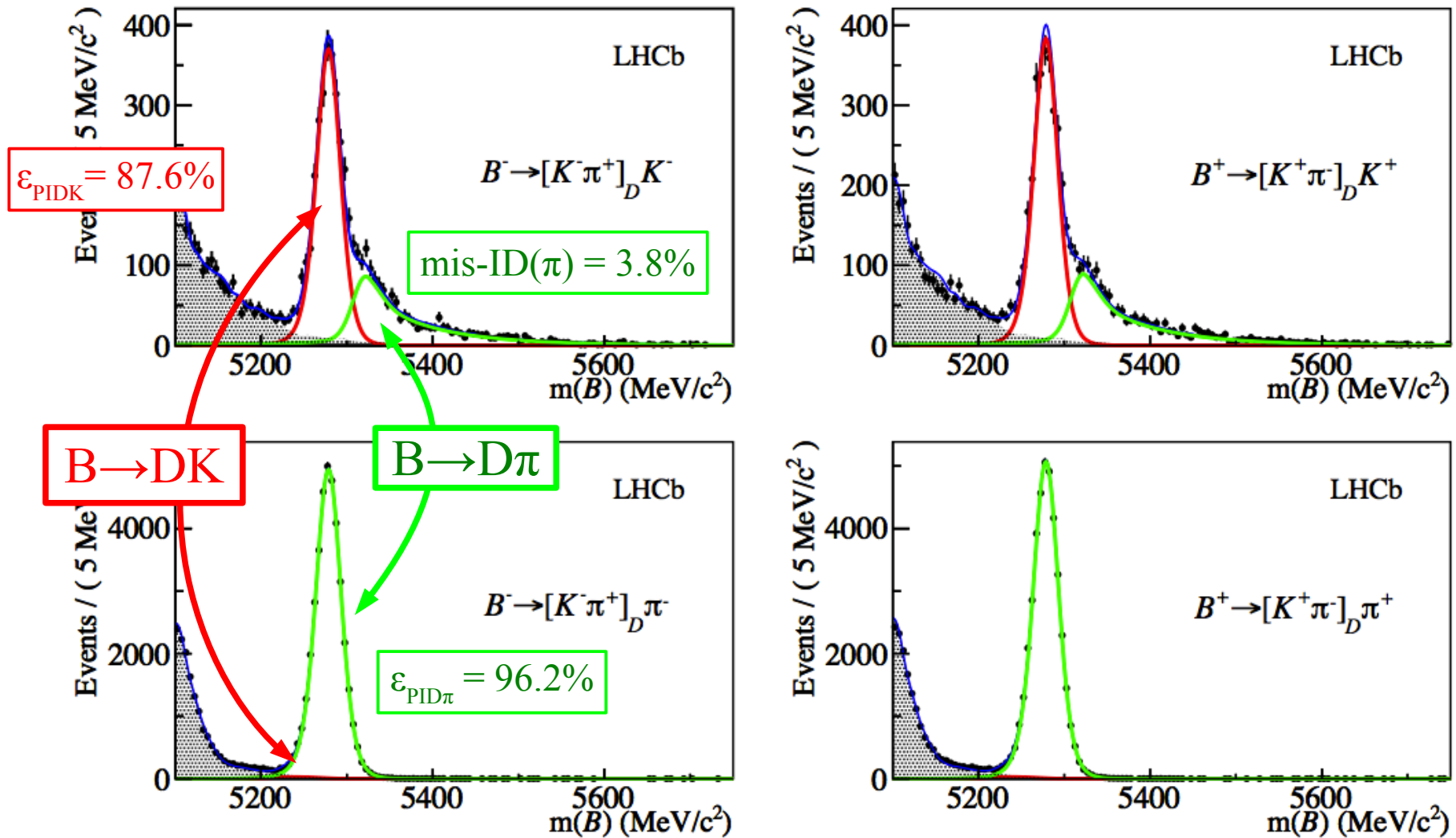
- Reconstructing neutral particles is difficult, so we miss the GLW CP- states.
- We also suffer in  $D \rightarrow K_S \pi \pi$ .
- At LHCb, we chose a more **“factory like” approach**.
- We use **all** reasonable ratios, including the asymmetry of the ADS favored modes.

$$\begin{aligned}
 R_{K/\pi}^{K\pi} &= 0.0774 \pm 0.0012 \pm 0.0018 \\
 R_{K/\pi}^{KK} &= 0.0773 \pm 0.0030 \pm 0.0018 \\
 R_{K/\pi}^{\pi\pi} &= 0.0803 \pm 0.0056 \pm 0.0017 \\
 A_{\pi}^{K\pi} &= -0.0001 \pm 0.0036 \pm 0.0095 \\
 A_K^{K\pi} &= 0.0044 \pm 0.0144 \pm 0.0174 \\
 A_K^{KK} &= 0.148 \pm 0.037 \pm 0.010 \\
 A_K^{\pi\pi} &= 0.135 \pm 0.066 \pm 0.010 \\
 A_{\pi}^{KK} &= -0.020 \pm 0.009 \pm 0.012 \\
 A_{\pi}^{\pi\pi} &= -0.001 \pm 0.017 \pm 0.010 \\
 R_K^- &= 0.0073 \pm 0.0023 \pm 0.0004 \\
 R_K^+ &= 0.0232 \pm 0.0034 \pm 0.0007 \\
 R_{\pi}^- &= 0.00469 \pm 0.00038 \pm 0.00008 \\
 R_{\pi}^+ &= 0.00352 \pm 0.00033 \pm 0.00007 .
 \end{aligned}$$

13 observables,  
all part of  
same analysis!

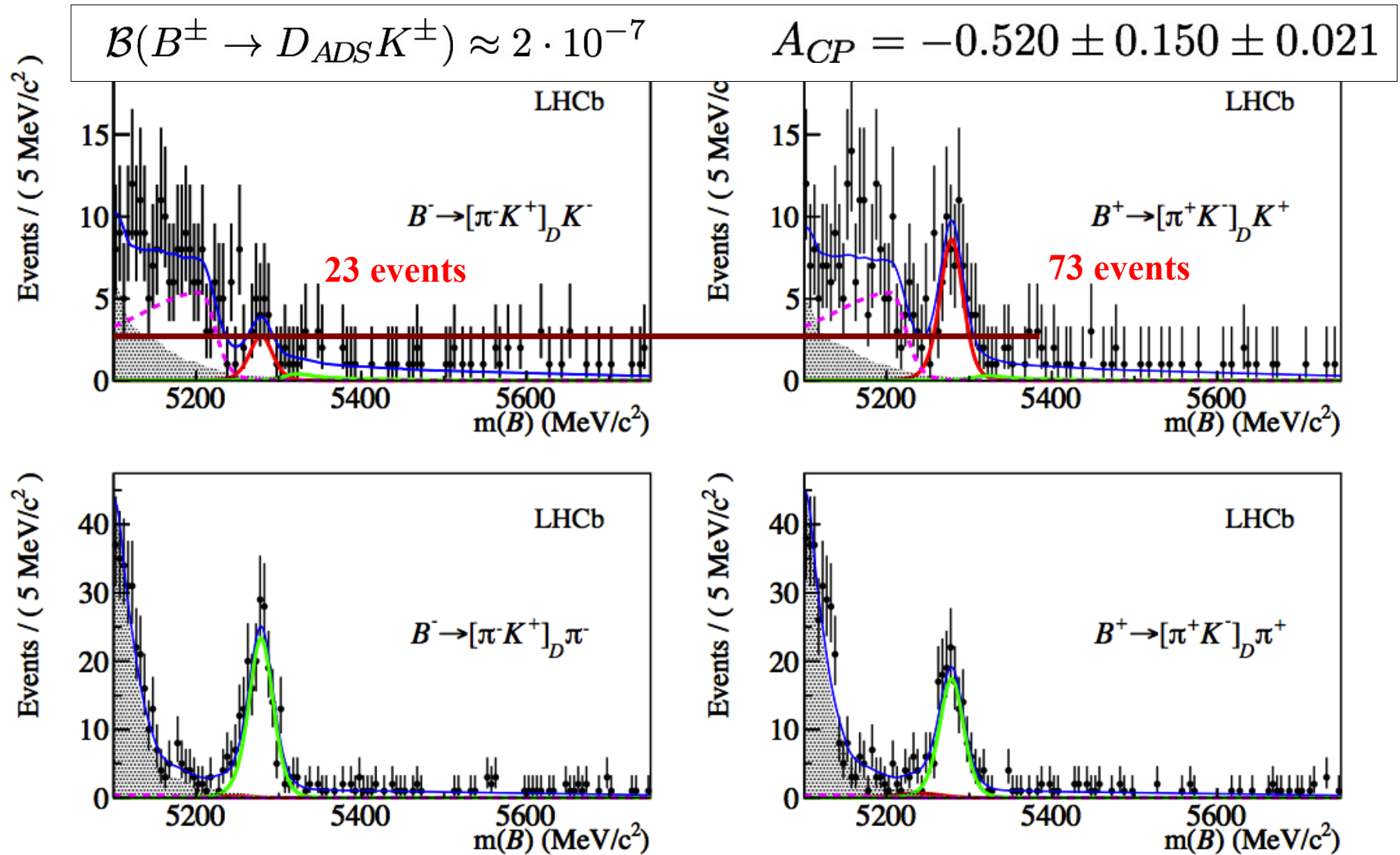
final states  $[f]_D$ :  
KK,  $\pi\pi$   
 $K\pi$ ,  $\pi K$

# $B \rightarrow D(K\pi)h$ : favored ADS mode



ARXIV:1203.3662

# $B \rightarrow D(\pi K)h$ : suppressed ADS mode



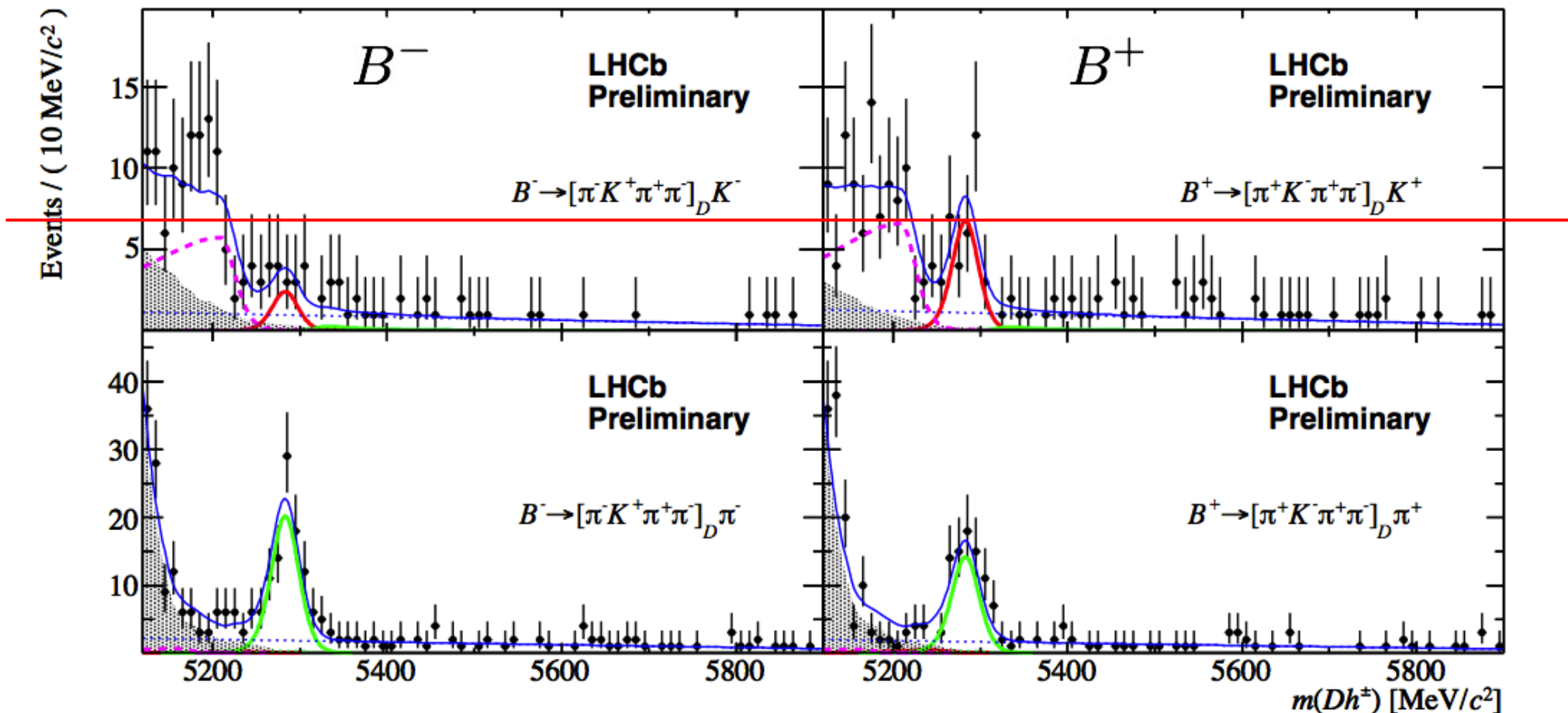
ARXIV:1203.3662

# four-body ADS

$B \rightarrow Dh$  followed by  $D \rightarrow K\pi\pi\pi$

First observations of these decay modes!

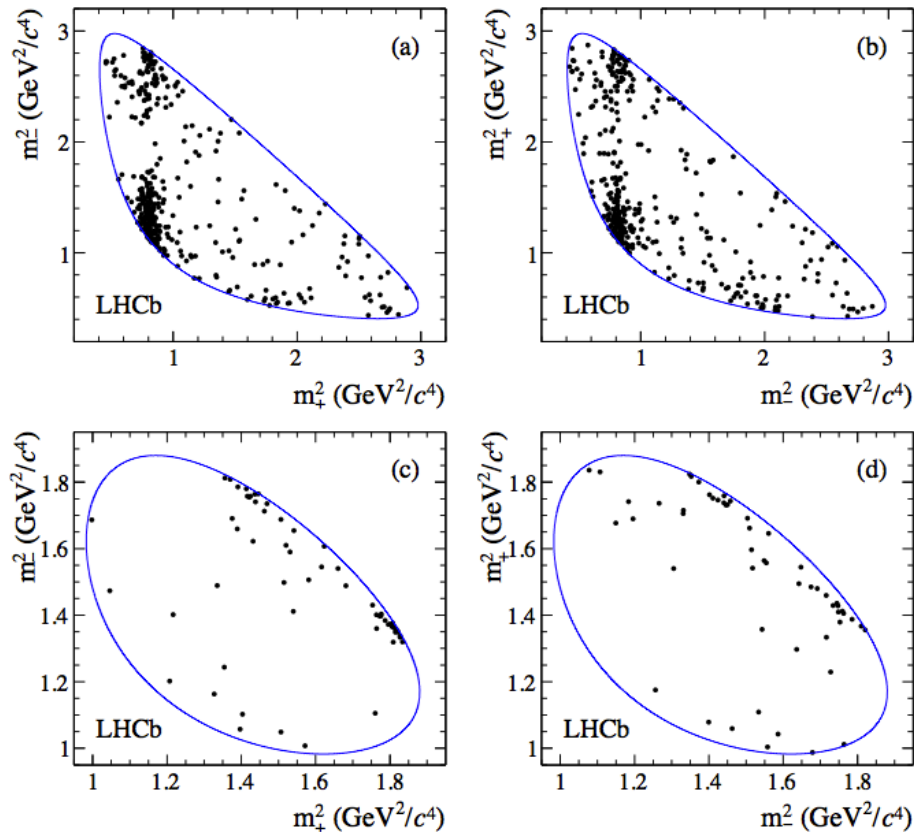
$$A_{CP} = -0.42 \pm 0.22$$





# model independent GGSZ

$$D^0 \rightarrow K_S^0 \pi^- \pi^+$$

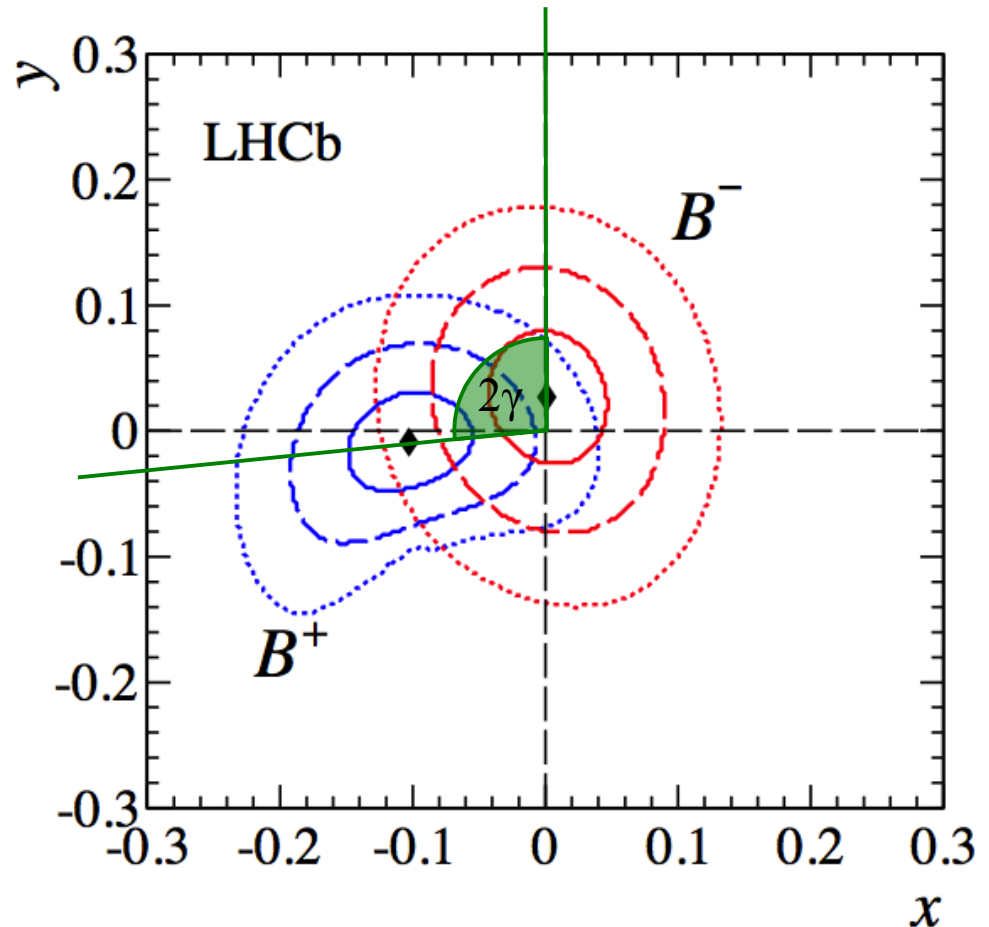


$$D^0 \rightarrow K_S^0 K^- K^+$$

Observables: “cartesian coordinates”

$$x_{\pm} = r_B \cos(\delta_B \pm \gamma)$$

$$y_{\pm} = r_B \sin(\delta_B \pm \gamma)$$



$$x_- = (0.0 \pm 4.3 \pm 1.5 \pm 0.6) \times 10^{-2}$$

$$y_- = (2.7 \pm 5.2 \pm 0.8 \pm 2.3) \times 10^{-2}$$

$$x_+ = (-10.3 \pm 4.5 \pm 1.8 \pm 1.4) \times 10^{-2}$$

$$y_+ = (-0.9 \pm 3.7 \pm 0.8 \pm 3.0) \times 10^{-2}$$

Part III:

gamma combination

# Combination

- We now have measured a total of **24  $\gamma$ -related observables**. What does it mean for  $\gamma$ ?
- At the CKM2012 conference, both Babar and Belle presented their legacy combinations.
- LHCb showed first results!
  - frequentist procedure
  - for the first time including the  $B \rightarrow D\pi$  system
  - considering CP violation in charm decays
  - partially considering charm mixing
- Combined likelihood:

$$\mathcal{L}(\vec{y}) = \frac{1}{N} \exp \left( -\frac{1}{2} (\vec{y} - \vec{y}_t)^T V_{\text{cov}}^{-1} (\vec{y} - \vec{y}_t) \right)$$
$$\chi^2(\vec{y}) = -2 \ln \mathcal{L}(\vec{y}) .$$

# statistical treatment

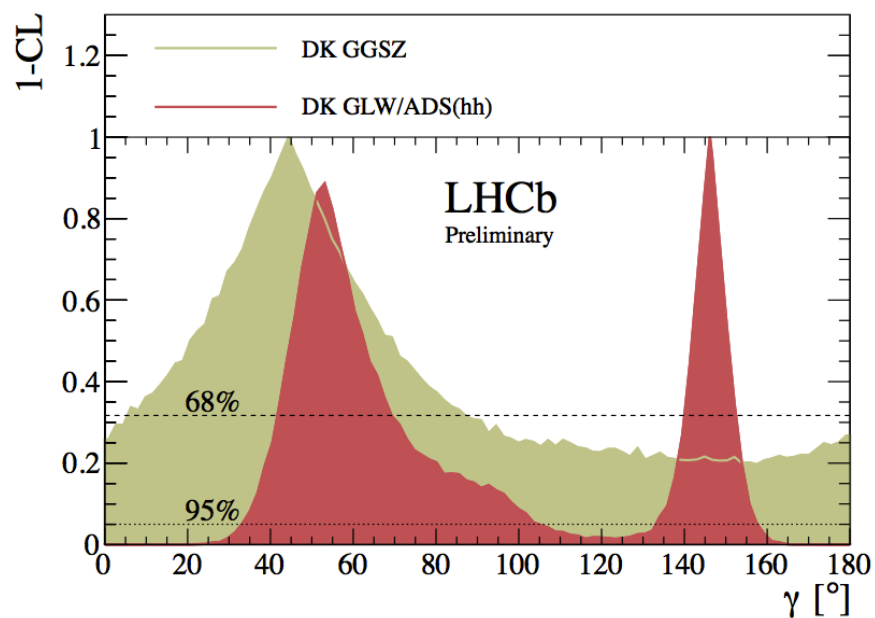
- The combined likelihood has a very rich structure:
  - many **nuisance parameters**
  - many trigonometrical functions, thus **many local minima**
  - **varying dimensionality** of the likelihood, if depends on the value of the nuisance parameters, potentially affecting the coverage
- Use a Feldman-Cousins based frequentist method (likelihood ratio ordering).
- Compute the actual distribution of the test statistic ( $\Delta\chi^2$ ) using toy Monte Carlo (“plug-in” method).
- **Test the coverage** at the best-fit point.

} direct product of  $r_B$  and angular terms:  
 $x_{\pm} = r_B \cos(\delta_B \pm \gamma)$

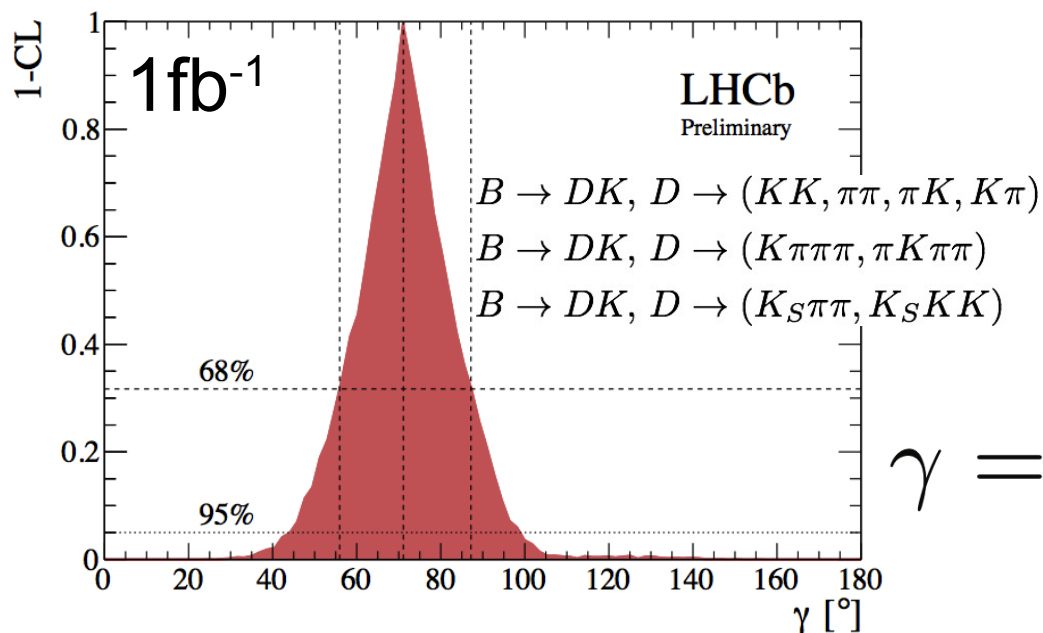
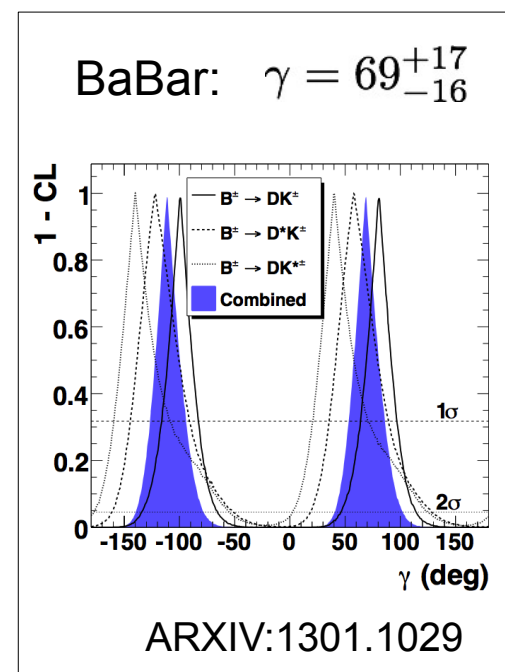
coverage test:

$D^0 K \& D^0 \pi$	$\alpha$
$\eta = 0.6827$	$0.6616 \pm 0.0067$
$\eta = 0.9545$	$0.9586 \pm 0.0028$
$\eta = 0.9973$	$0.9958 \pm 0.0009$

# B → DK only

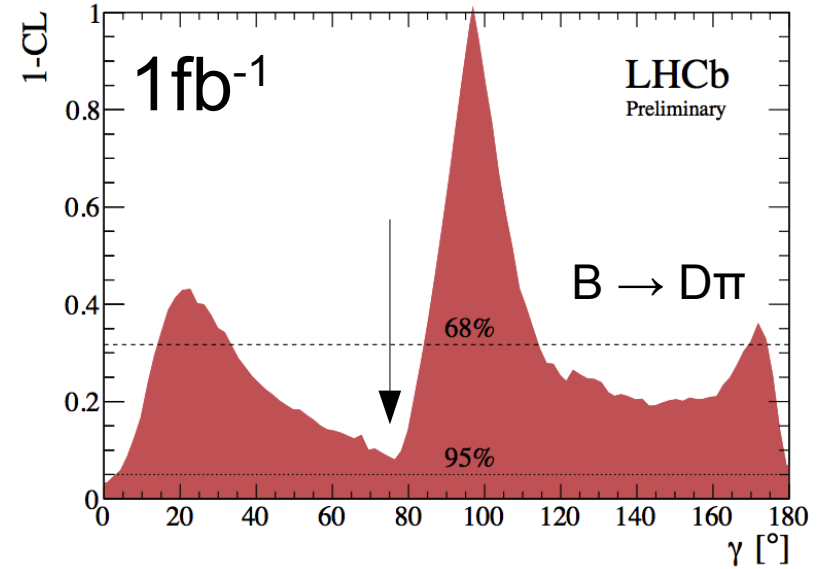
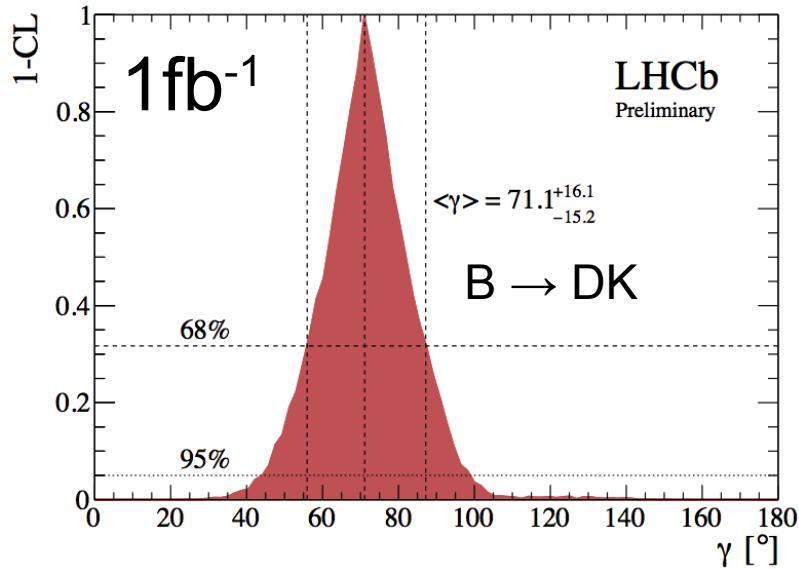


This is about as precise as the B-factory legacy combinations!

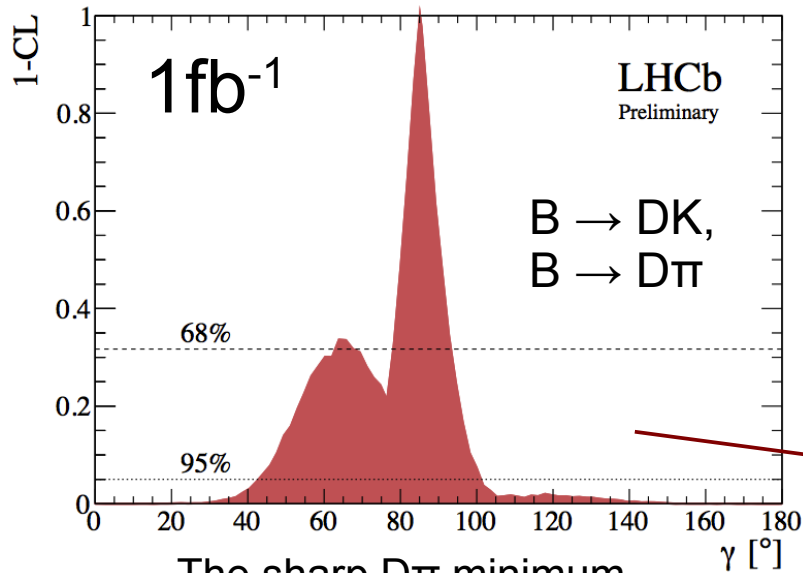


$$\gamma = (71.1^{+16.6}_{-15.7})^\circ$$

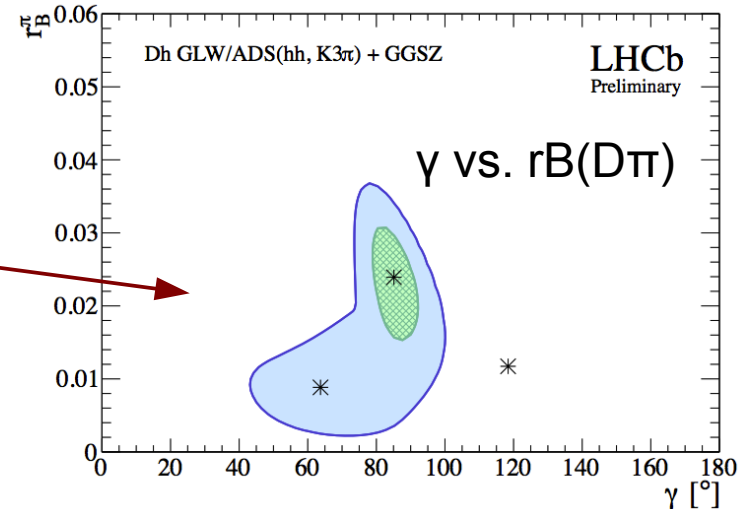
# adding $B \rightarrow D\pi$



$B \rightarrow D\pi$  is more precise than expected.  
But (almost) no constraint at  $2\sigma$ .



The sharp  $D\pi$  minimum affects the full combination!

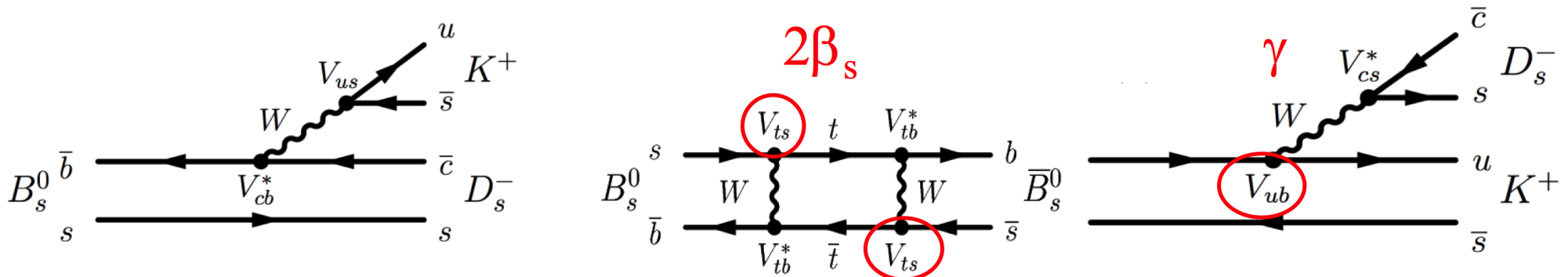


Part IV:

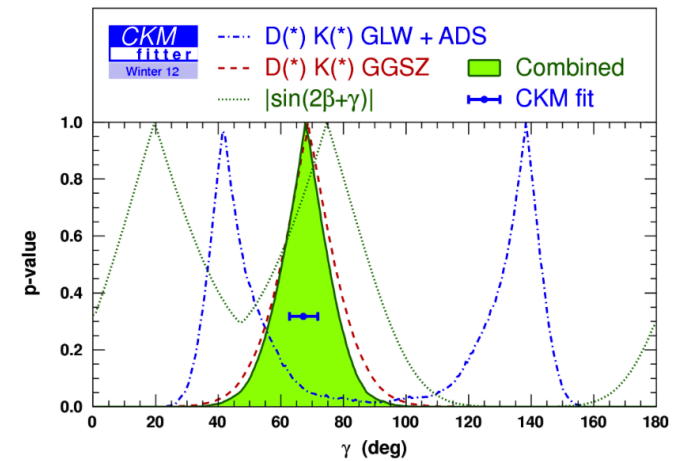
time dependent measurements

# time dependent measurements

- It is also possible to use tree decays of neutral B mesons [1]!
- Using charged-particle final states, interference is achieved through mixing.



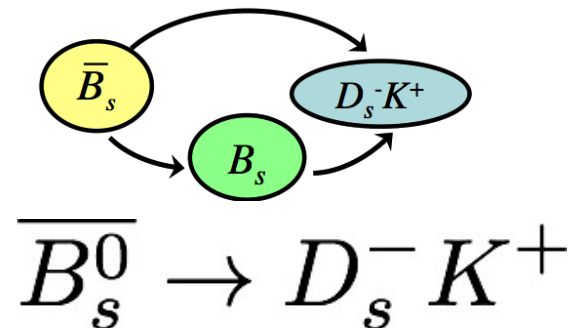
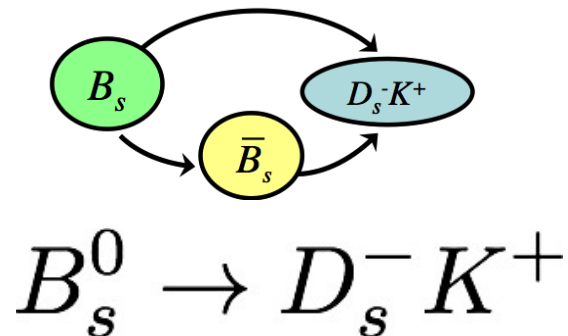
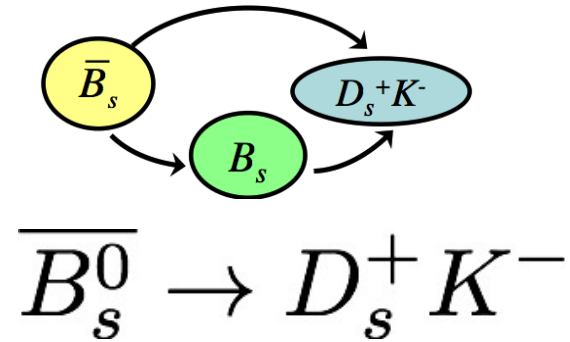
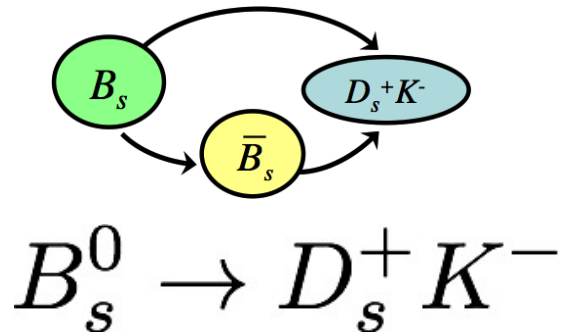
- B-factories performed such measurements with  $B^0 \rightarrow D^+ \pi^-$ , constraining  $\sin(2\beta+\gamma)$
- **Much better sensitivity expected in  $B_s \rightarrow D_s K$ ,** comparable to GGSZ, only really possible at LHCb
  - large amplitude ratio:  $r_{D_s K} \sim 0.4$
  - finite decay width difference:  
 $\Delta\Gamma = 0.091 \pm 0.011 \text{ ps}^{-1}$  (HFAG fall 2012)



[1] R. Fleischer. New strategies to obtain insights into CP violation through  $B_{(s)} \rightarrow D_{(s)}^\pm K^\mp, D_{(s)}^{*\pm} K^\mp, \dots$  and  $B_{(d)} \rightarrow D^\pm \pi^\mp, D^{*\pm} \pi^\mp, \dots$  decays. *Nucl.Phys.*, B671:459–482, 2003.



# four decay rates



each has their own time dependence

# four decay rates

$$\frac{d\Gamma_{B_s^0 \rightarrow f}(t)}{dt e^{-\Gamma_s t}} = \frac{1}{2} |A_f|^2 (1 + |\lambda_f|^2) \left[ \cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + D_f \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) + C_f \cos(\Delta m_s t) - S_f \sin(\Delta m_s t) \right] \quad (1)$$

$$\frac{d\Gamma_{\bar{B}_s^0 \rightarrow f}(t)}{dt e^{-\Gamma_s t}} = \frac{1}{2} |A_f|^2 \left| \frac{p}{q} \right|^2 (1 + |\lambda_f|^2) \left[ \cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + D_f \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) - C_f \cos(\Delta m_s t) + S_f \sin(\Delta m_s t) \right] \quad (2)$$

$$\frac{d\Gamma_{\bar{B}_s^0 \rightarrow \bar{f}}(t)}{dt e^{-\Gamma_s t}} = \frac{1}{2} |\bar{A}_{\bar{f}}|^2 (1 + |\bar{\lambda}_{\bar{f}}|^2) \left[ \cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + D_{\bar{f}} \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) + C_{\bar{f}} \cos(\Delta m_s t) - S_{\bar{f}} \sin(\Delta m_s t) \right] \quad (3)$$

$$\frac{d\Gamma_{B_s^0 \rightarrow \bar{f}}(t)}{dt e^{-\Gamma_s t}} = \frac{1}{2} |\bar{A}_{\bar{f}}|^2 \left| \frac{q}{p} \right|^2 (1 + |\bar{\lambda}_{\bar{f}}|^2) \left[ \cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + D_{\bar{f}} \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) - C_{\bar{f}} \cos(\Delta m_s t) + S_{\bar{f}} \sin(\Delta m_s t) \right] \quad (4)$$

# observables

Five observables:

$$C = \frac{1 - r_{D_s K}^2}{1 + r_{D_s K}^2},$$

$$D_f = \frac{2r_{D_s K} \cos(\Delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}, \quad D_{\bar{f}} = \frac{2r_{D_s K} \cos(\Delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2},$$

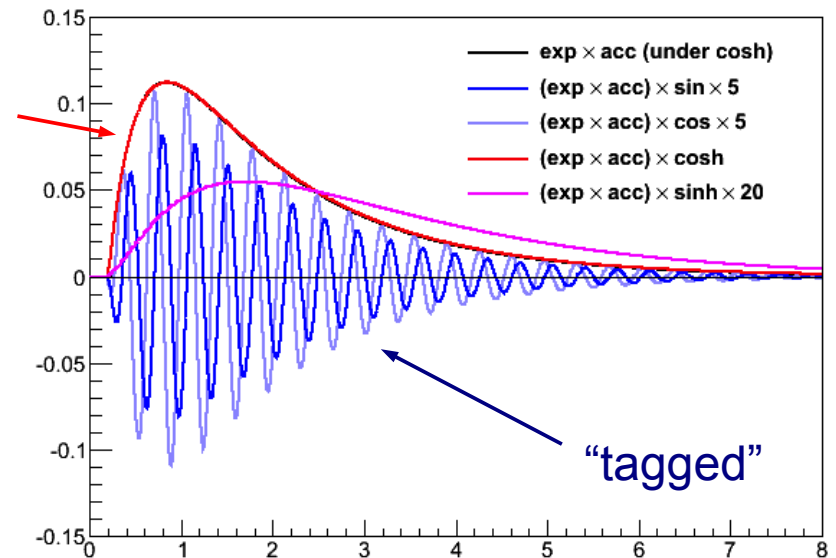
$$S_f = \frac{2r_{D_s K} \sin(\Delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}, \quad S_{\bar{f}} = \frac{2r_{D_s K} \sin(\Delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}.$$

Measurement:

- excellent time resolution needed to resolve fast  $B_s$  oscillations
- flavor tagging to tell  $B_s$  from  $\overline{B}_s$

$$\epsilon D^2 = 1.9\%$$

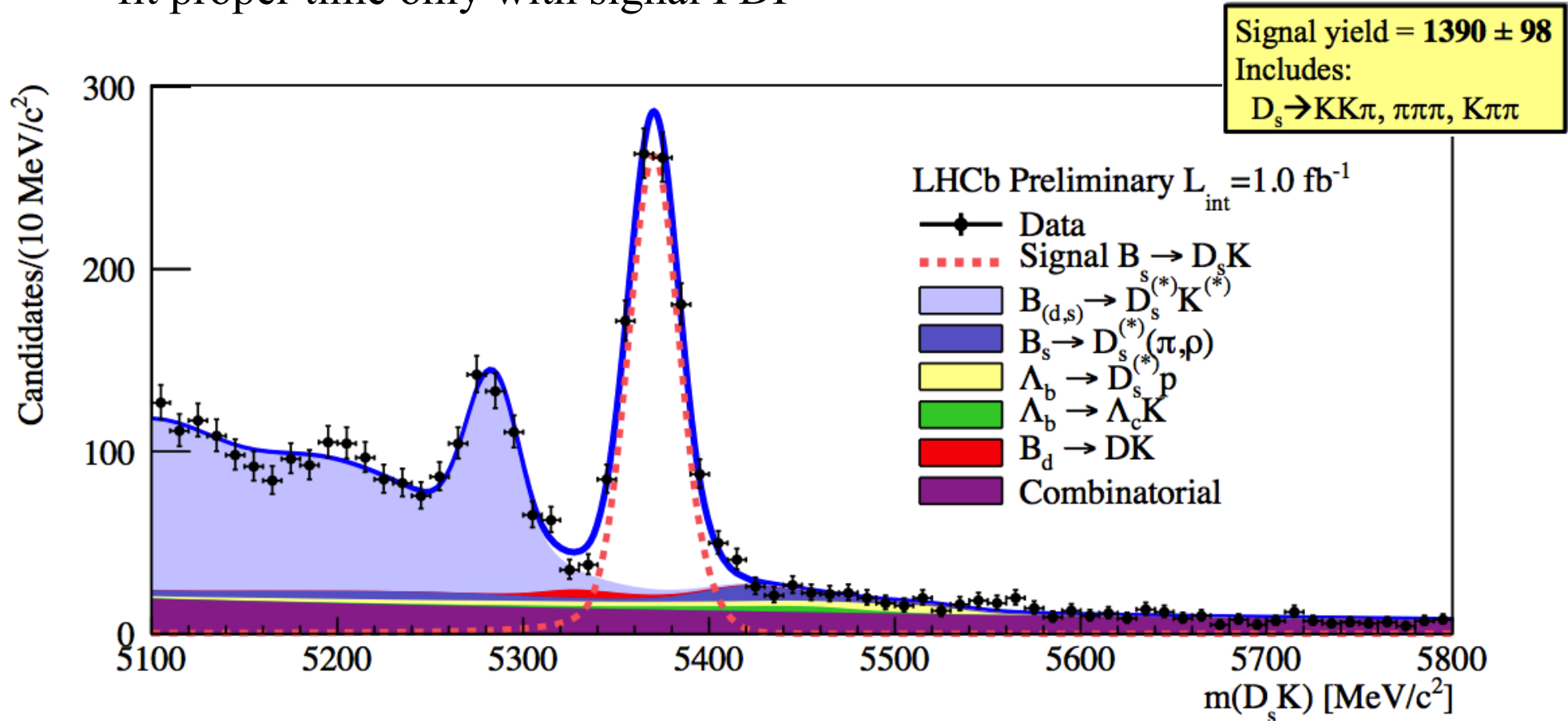
“untagged”



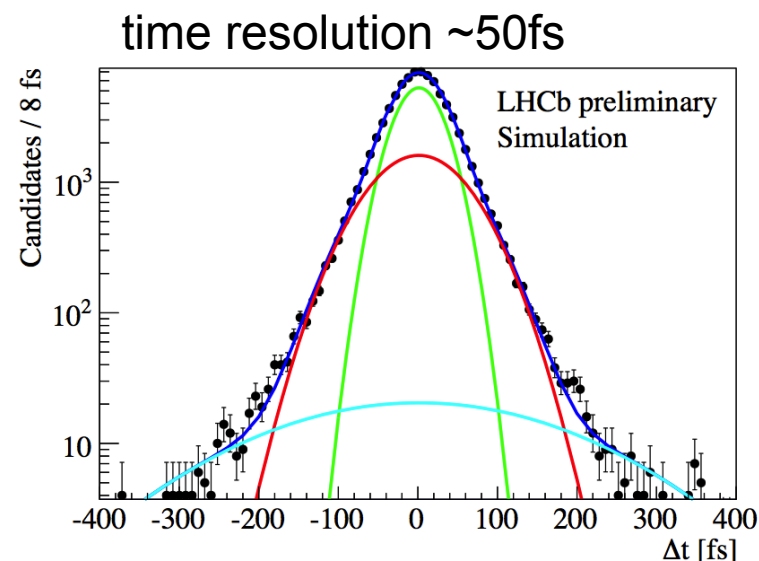
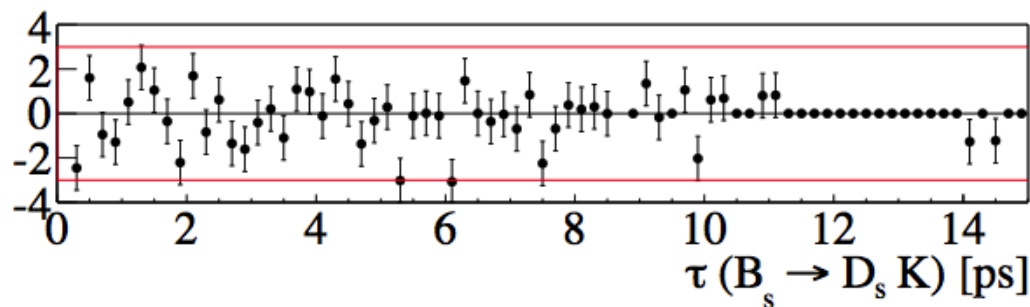
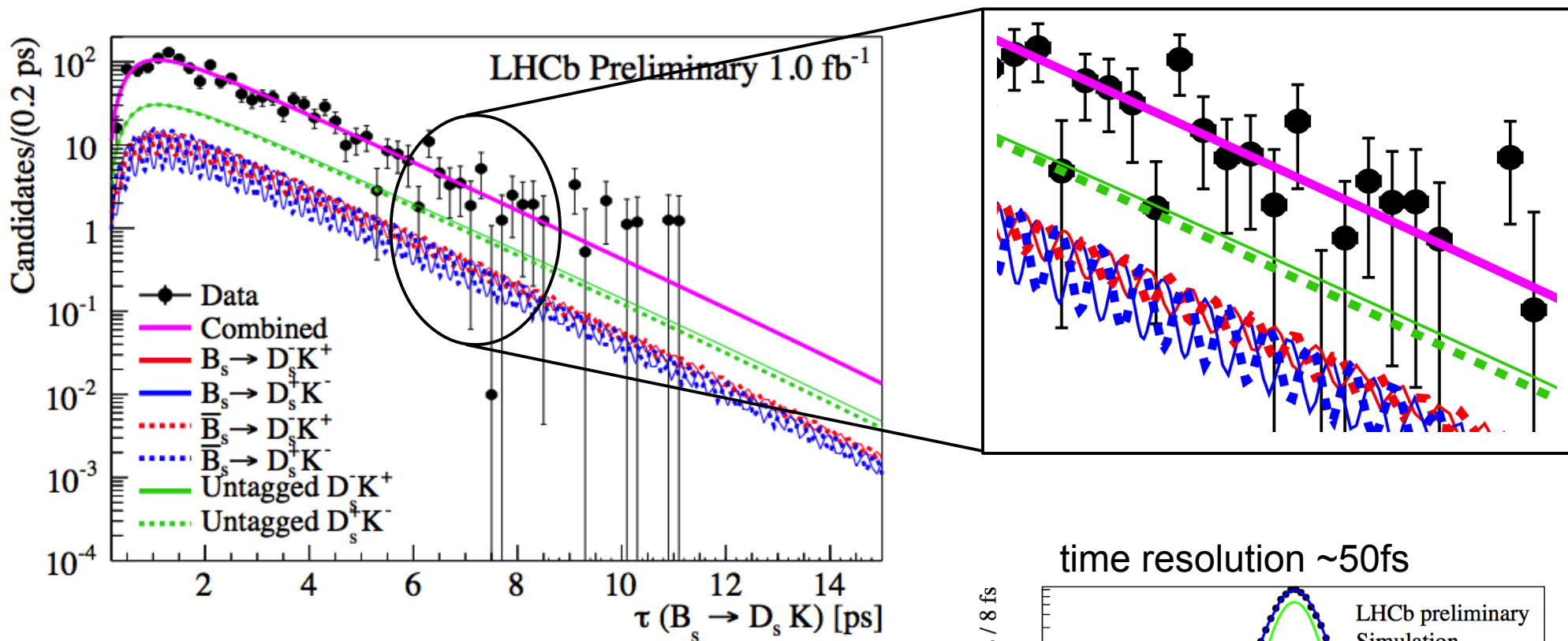
time [ps]

# mass plot

- rich physics backgrounds – all could have different time structure!
- employ an “sFit” technique (“cFit” available):
  - use  $B_s$  mass as discriminating variable to compute per-event weights
  - fit proper time only with signal PDF



# time dependent DsK



# result

	$C$	$S_f$	$S_{\bar{f}}$	$D_f$	$D_{\bar{f}}$
Toy corrected central value	1.01	-1.25	0.08	-1.33	-0.81
Statistical uncertainty	0.50	0.56	0.68	0.60	0.56
Systematic uncertainties ( $\sigma_{\text{stat}}$ )					
Decay-time bias	0.03	0.05	0.05	0.00	0.00
Decay-time resolution	0.11	0.08	0.09	0.00	0.00
Tagging calibration	0.23	0.17	0.16	0.00	0.00
Backgrounds	0.15	0.07	0.07	0.07	0.07
Fixed parameters	0.15	0.22	0.20	0.40	0.42
Asymmetries	0.12	0.01	0.04	0.00	0.02
Momentum/length scale	0.00	0.00	0.00	0.00	0.00
k-factors	0.27	0.27	0.27	0.08	0.08
Bias correction	0.03	0.03	0.03	0.03	0.03
Total systematic ( $\sigma_{\text{stat}}$ )	0.46	0.50	0.35	0.43	0.46

- statistically limited, but non-negligible systematic errors
- no constraints on  $\gamma$  yet (warning: syst. correlations cannot be neglected)
- significant update planned

# Conclusion

# Conclusion

- In the era of precision flavor physics,  $\gamma$  provides a set point of the Standard Model because of its **negligible theoretical uncertainty** (tree decays).
  - does the unitarity triangle close?
  - can probe for new physics at very high energy scales
- First LHCb measurements are available.
  - been statistically lucky with  $B \rightarrow D\pi$
- with  $1\text{fb}^{-1}$  already at the precision of the B factories
- For the first time, attempt to measure  $\gamma$  in time-dependent analysis of  $B_s \rightarrow D_s K$  - high precision seems possible!

$$\gamma = (71.1^{+16.6}_{-15.7})^\circ \quad B \rightarrow DK \text{ only, LHCb at CKM2012 preliminary}$$



# Backup

# flavor tagging

